Quantum dots in three dimensions

FYS3140 – Project 2

[github.com/henrikx2/FYS3150](https://github.com/henrikx2/FYS3150)

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# Abstract

In this report, the eigenvalue problem of a system with electrons in a harmonic oscillator potential is explored. The well-known and stable Jacobi’s rotation algorithm is used to diagonalize the Hamiltonian matrix into an eigenvalue matrix. The equilibrium between calculation speed and decimal precision of the eigenvalues are discussed as a function of step size, interval length and numbers of iterations.

# Introduction

The experiment aims to solve the general eigenvalue equation of the type

Where is going to be the Hamiltonian operator of the harmonic oscillator well potential and is the energy eigenvalue of the given state .

The equation is solved by firstly; constructing a tridiagonal Hamiltonian matrix. This matrix is derived by simplifying the Schroedinger’s equation for one or two electrons in the HO potential and making it dimensionless. Then, the matrix is processed with orthogonal transformation through Jacobi’s rotation algorithm. This reveals a diagonal matrix with the eigenvalues of the Hamiltonian at its diagonal. The eigenvalues are then compared to the analytic eigenvalues. Lastly, the ground state of the system’s dependence on Coulomb interactions and varying HO potential will be explored.

# Theory

## One electron in the HO potential

The solution to the radial part of the Schroedinger’s equation for one electron in the HO potential is given as

|  |  |  |
| --- | --- | --- |
|  |  | (1) |

Where is the HO potential with and , is the wavefunction and is the energy of the harmonic oscillator. The energies can be written as a function of the frequency as

|  |  |  |
| --- | --- | --- |
|  |  | (2) |

Where is the principal quantum number and is the angular momentum quantum number of the electron. Substituting , Eq. 1 becomes

|  |  |  |
| --- | --- | --- |
|  |  | (3) |

Where the Dirichlet boundary conditions give and . To make the Eq. 3 dimensionless, it is rewritten as a function of the dimensionless variable where has dimensions length. The HO potential is then given . This project only concentrates only on a system where , which simplifies Eq. 3 to

|  |  |  |
| --- | --- | --- |
|  |  | (4) |

To make Eq. 4 suitable for solving numerically, it is multiplied by , so that

|  |  |  |
| --- | --- | --- |
|  |  | (5) |

Which is further manipulated with the condition

|  |  |  |
| --- | --- | --- |
|  |  | (6) |

So that defining

|  |  |  |
| --- | --- | --- |
|  |  | (7) |

Gives the simplified Schroedinger’s equation which is to be solved as

|  |  |  |
| --- | --- | --- |
|  |  | (8) |

The analytical eigenvalues of this equation are

Since the maximum value cannot be set to , this will be varying as to get the most correct energies. The values of is defined . The number of mesh points are given as and the step length reads

|  |  |  |
| --- | --- | --- |
|  |  | (9) |

Which in turn gives the expression for the value of at a given iteration as

|  |  |  |
| --- | --- | --- |
|  |  | (10) |

An approximation to the 2nd derivative can be expressed

|  |  |  |
| --- | --- | --- |
|  |  | (11) |

Using Eq. 11 and the HO potential to describe the discrete case of the Schroedinger’s equation as

|  |  |  |
| --- | --- | --- |
|  |  | (12) |

It is easy to notice the tridiagonal matrix equation where the matrix is defined as follows

In other words, the diagonal elements and the non-diagonal elements are denoted as

|  |  |  |
| --- | --- | --- |
|  |  | (13) |
|  |  | (14) | |

## Two electrons in the HO potential

Now, the two-electron system is a bit more exciting, because it opens the door to study the coulomb interactions between electrons. The general Schroedinger’s equation for one electron in a HO potential well is as before

|  |  |  |
| --- | --- | --- |
|  |  | (15) |

Where is the energy of the system with only one electron. Now, adding the other electron (but no Coulomb interactions) the equation reads

|  |  |  |
| --- | --- | --- |
|  |  | (16) |

Now, using relative coordinate and center-of-mass coordinate , the Shroedinger equation becomes

|  |  |  |
| --- | --- | --- |
|  |  | (17) |

Separating and can be done using the ansatz for the separation of the wave function , where det energy is a sum of the relative and the center-of-mass energy given by .

The Coulomb interactions between the electrons is given by the expression

|  |  |  |
| --- | --- | --- |
|  |  | (18) |

In which the constant . This can be added to the part of the equation

|  |  |  |
| --- | --- | --- |
|  |  | (19) |

By making Eq. 19 dimensionless by the same steps as in 2.1, the equation reads

|  |  |  |
| --- | --- | --- |
|  |  | (20) |

To make Eq. 20 “identical” to Eq. 8, a new “frequency” is defined as

|  |  |  |
| --- | --- | --- |
|  |  | (21) |

Where the parameter mirrors the strength of the HO potential. Furthermore

|  |  |  |
| --- | --- | --- |
|  |  | (22) |

Such that

|  |  |  |
| --- | --- | --- |
|  |  | (23) |

And the eigenvalue

|  |  |  |
| --- | --- | --- |
|  |  | (24) |

So that the final Schroedinger’s equation reads

|  |  |  |
| --- | --- | --- |
|  |  | (25) |

This equation has the analytical eigenvalues given by Eq. 26[[1]](#footnote-1) underneath:

|  |  |  |
| --- | --- | --- |
|  |  | (26) |

Where:

The expression for the analytical eigenvalues with no electron-electron interaction, is given as:

|  |  |  |
| --- | --- | --- |
|  |  | (27) |

These expressions will be used to evaluate the experimental results from the program later.

The diagonal in the tridiagonal matrix is now defined (the off diagonals stay the same)

|  |  |  |
| --- | --- | --- |
|  |  | (28) |

If there where not to be a Coulomb interaction between the two particles, the would be just like the one in the singe electron system, but with and hence would be defined as:

## Preservation of orthogonality

For the method described in section 3 to work, it is essential to prove that a unitary transformation preserves the dot product and the orthogonality of vectors in a basis. Consider the orthogonal basis

For a matrix U to be unitary it must have the property . Consider the transformation . The dot product is then

Hence, the orthogonal/unitary transformation preserves the orthogonality◾

# Methods

## Jacobi’s rotation algorithm

Jacobi’s rotation algorithm evolves around performing an orthogonal transformation with a rotation matrix of the type

Which has the property and performs a rotation of angle in the Euclidean space. The elements of that are not zero are given as

The algorithm is used to transform a symmetric matrix into a diagonal matrix where the diagonal consists of the eigenvalues. This diagonalization is performed by doing the orthogonal transformation several times, so that the off-diagonal elements converge to zero.

Defining and a 2x2 system would look like this

The requirement that implies that

Where it’s easy to see that if the off-diagonal element , then , which implies that and , or in other words; no rotation needed as off-diagonals is close to zero. Next, defining

A quadratic equation can be derived from the relation so that

|  |  |  |
| --- | --- | --- |
|  |  | (29) |

From the relation , and are defined as

The equations for the new matrix elements when iterating are:

This gives a new matrix , on which the same procedure is done until it’s off-diagonal elements are under a certain tolerance

The rate of which the Jacobi’s method converges to the diagonal matrix with the eigenvalues is horribly slow. It takes somewhat rotations to make the off-diagonals zero, and each rotation uses approximately FLOPS. This makes the algorithm use at least somewhere around FLOPS.

## Code structure

1. Build the matrix A by implementing the diagonal elements from the corresponding problem in section 2.
2. Find the off-diagonal element with the largest absolute value and assign its matrix index to and .
3. Calculate , and .
4. Calculate the 6 equations with the obtained and and update the matrix A.
5. Repeat until off-diagonal elements are lower than the specified tolerance.

## Running the calculations

To run the calculations using the provided source code in the git repository, it is required to have Python 3 with the NumPy, SciPy, matplotlib and numba packages installed. Run the calculations by executing main.py, this will start a UI. Choose to ether; calculate all necessary data and make plots, just plot data which is already calculated or do individual calculation with free variables.

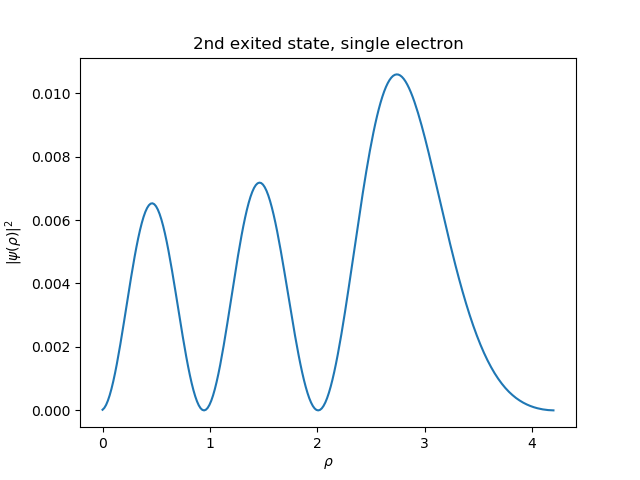
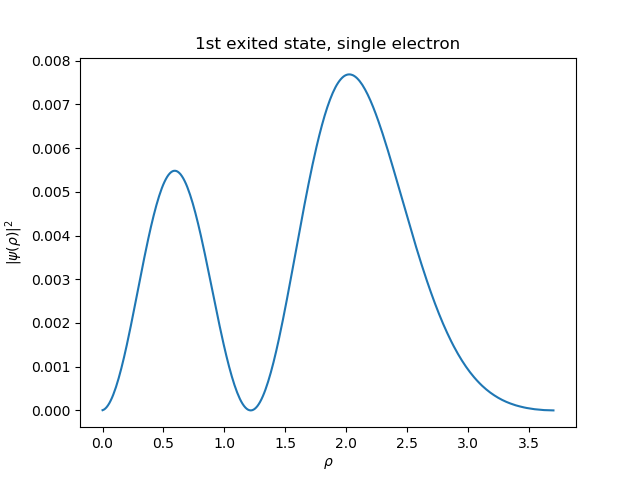
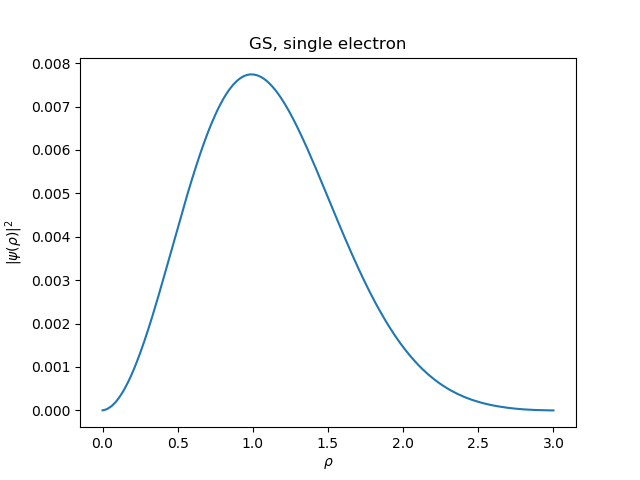
# Results and discussion

## Single electron

The single electron system has the analytical eigenvalues . To achieve these eigenvalues, it was necessary to run the algorithm several times trying different values of and . It was also found that it was not possible to achieve 4 decimal precision of the eigenvalues of the third lowest states at the same time. This resulted in Table 6.2.1, which shows which conditions on and that gave the eigenvalues with 4 decimal precision.

Figure 4.1.1a-c shows how the probability distribution, , of the three lowest states look like with the conditions from Table 6.2.1. From these figures, it is easy to notice how the higher excited states have a larger spatial distribution than that of the lower energetic states.

The limitations on how many eigenvalues that can be calculated at the same time evolves from the fact that infinity has to be assigned to something smaller than infinity. This makes it possible to get one eigenvalue correct (using a satisfying value of and number of integration points), but the other values will likely be far off. To achieve several correct eigenvalues in the same matrix, the matrix would have to be very large (high ) and with a large value of .



(c))

(b)

(a))

Figure 4.1.1: The probability density of the wave function in ground state (a), 1st exited state (b) and 2nd exited state (c). Calculated using the experimental eigenvalues with 4 decimal precision.

## Two electrons

In the two-electron system it is exciting to study how the Coulomb interactions impact how different values of impact the wavefunction and the energies of the ground state. So, first off is plotting the two different systems for values , , and . One should be able to predict that the larger value of ; the smaller the distribution of the wavefunction. This is shown in Figure 4.2.1 and Figure 4.2.2.

In these figures, it is clear how the HO frequency affects the wavefunction of the system to broaden in -space. This is intuitive, because a smaller potential would give the electrons more possibility of existing in a larger area.

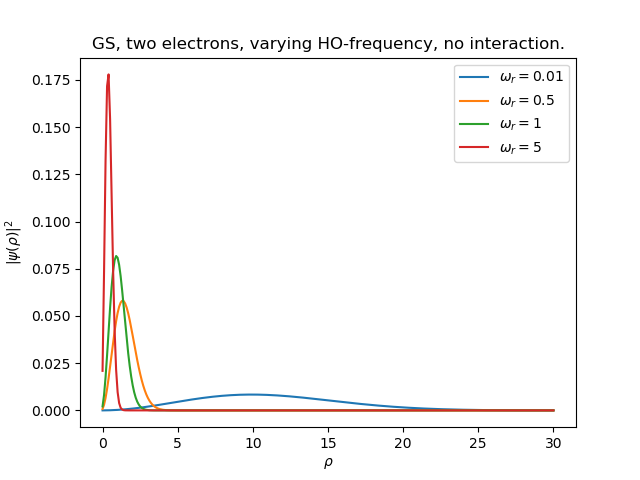


Figure 4.2.1: The probability density of the GS wavefunction without the electron-electron interaction for different values of the HO-potential strength . See how the distribution broadens for smaller potentials.

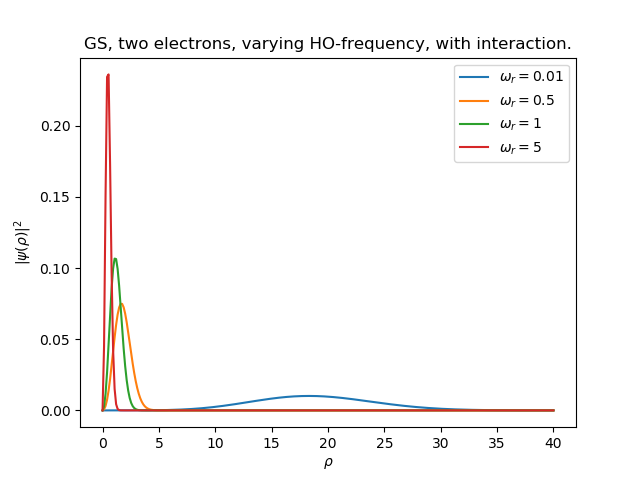


Figure 4.2.2: The probability density of the GS wavefunction with the electron-electron interaction for different values of the HO-potential strength . See how the distribution broadens for smaller potentials.

To figure out how the Coulomb interactions affect the systems energy, one may look at how the peaks in Figure 4.2.1 and Figure 4.2.2 are shifted. But this is hard to see with the naked eye, especially for the most energetic states. To see the difference between the two systems, let’s look at the experimental eigenvalues. In Table 4.2.1, the eigenvalues in the two cases are compared as function of . The value of and are manipulated so that the wavefunction with has the most correct eigenvalue (such that the error in the eigenvalues increases as increases). Table 6.1.1 in Appendix lists the analytical eigenvalues of the two-electron system with and without the Coulombs interactions.

*Table 4.2.1: Comparison of eigenvalues in the Coulomb interacting systems and the non-Coulomb interacting systems.*

|  |  |  |  |
| --- | --- | --- | --- |
|  | **Two electrons, no interaction** | **Two electrons, with interaction** | **Increase in energy from interaction:** |
| **HO frequency,** | **Experimental eigenvalues, :** | **Experimental**  **eigenvalues, :** |
| 0.01 | 0.029990 | 0.105876 | 0.075886 |
| 0.5 | 1.426043 | 2.232772 | 0.806729 |
| 1 | 2.788140 | 4.060586 | 1.272446 |
| 5 | 12.642609 | 17.346902 | 4.704293 |

So, the energy of the system with Coulomb interaction has of course an increase in energy because of the repulsion forces between the electrons. By looking at how much the increase in energy is due to the Coulomb interaction as function of the HO-potential strength , one can see that the relation is at least not linear. This is likely since the higher the is, the narrower the peak of the probability density is (Figure 4.2.1 and Figure 4.2.2 shows this), and therefore the closer the electrons must be. Which of course gives a large repulsion potential. As mentioned earlier, from Figure 4.2.1 and 4.2.2, one can also see how the Coulomb interactions actually makes the squared wave function broader.

## Jacobi’s method

As much as this is a report solving the Schroedinger’s equation for 1-2 electrons in a harmonic oscillator potential and analysing the aspects of this, there is also a numerical algorithm to discuss. One interesting aspect to investigate is how many similarity transformations are needed as function of the dimension of the matrix . Figure 4.2.1 shows a plot of exactly this.

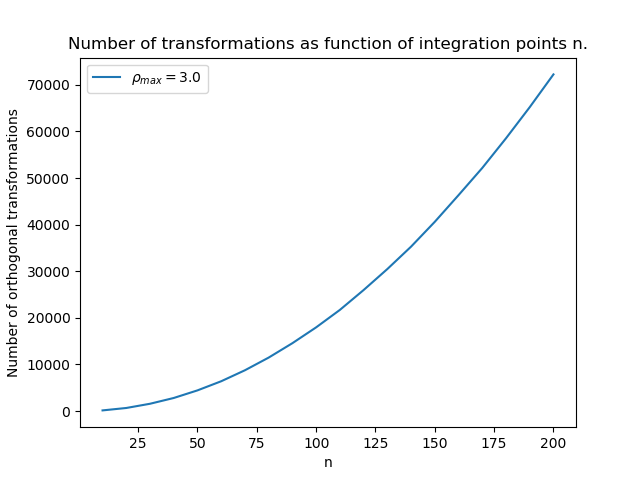


Figure 4.3.1: Number of similarity transformations as function of number of integration points in the Jacobi’s algorithm. This should in theory be a 2nd degree curve.

Using the SciPy function curve\_fit to interpret the points in this curve as a 2nd degree polynomial gives the equation:

This estimation of the number of similarity transformations falls somewhat short of the theoretic number () of rotations, but perhaps a larger interval of could make the function more like the theoretic assumptions. When it comes to speed, Jacobi’s method is a slow, but stable algorithm. One of the limiting factors when it comes to time usage are the integration points, . As increases, the double for-loop that finds the maximum off-diagonal element is rapidly getting more time consuming.

Though the algorithm itself is rather time consuming, it has no problems computing the eigenvalues of symmetric matrices of sizes up to (with a tolerance where ), if one has the time to wait a few minutes. The actual computation is in fact very precise, of course depending on the toleration limit for the zeros on the diagonal of the transformed matrix. This is also what gives the error of the method.

# Conclusion

# Appendix

## Analytical eigenvalues

Using Eq. 26 and 27 with the fact that and , the eigenvalues in Table 6.1.1 may be calculated.

*Table 6.1.1: The analytical eigenvalues of the two-electron system with different frequencies.*

|  |  |  |
| --- | --- | --- |
|  | **Two electrons, no interaction** | **Two electrons, with interaction** |
| **HO frequency,** | **Analytical eigenvalues, :** | **Analytical eigenvalues, :** |
| 0.01 | 0.03 | 0.1050 |
| 0.5 | 1.50 | 2.0566 |
| 1 | 3.00 | 3.6220 |
| 5 | 15.00 | 14.1864 |

## Optimal values of and

Table 6.2.1 shows the optimal values of and to achieve 4 decimal precision of the three lowest energy eigenvalues of the single electron system.

*Table 6.2.1: Optimal values of and to achieve 4 decimal precision of eigenvalues in the one-electron system.*

|  |  |  |  |
| --- | --- | --- | --- |
| **Analytical eigenvalues** |  |  | **Experimental eigenvalues** |
| 3 | 3 | 325 | 2.999987 |
| 7 | 3.7 | 337 | 7.000073 |
| 11 | 4.2 | 254 | 10.999960 |

# References

1. **Hjorth-Jensen, M. (2019)**, *Computational Physics, Project 2 Fall 2019*, Department of Physics, University of Oslo. <https://github.com/CompPhysics/ComputationalPhysics/blob/master/doc/Projects/2019/Project2/pdf/Project2.pdf>
2. **Hjorth-Jensen, M. (2019)**, *Computational Physics, Eigenvalue Problems,* Department of Physics, University of Oslo. <http://compphysics.github.io/ComputationalPhysics/doc/pub/eigvalues/html/eigvalues.html>
3. Analytical values <https://journals.aps.org/pra/pdf/10.1103/PhysRevA.48.3561>

1. *Eq. 26 and 27 are originating from equation (16a) and (16b) in the article: M. Taut, Phys. Rev. A 48, 3561 (1993).* [↑](#footnote-ref-1)